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## Characterization of chemical reaction on heat transfer through the nano fluid

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### Abstract

The chemical reaction of metal particle with the solvent is natural and decomposition of the metal particle is a natural phenomenon. But, this decomposition affects the viscosity and the fluid flow. The inclusion of chemical reaction rate parameter in the expression of viscosity given by Graham A.L (Applied Science Research, 37(3), pp.275-286(1981)) address the behavior of the viscosity of the Nano fluid in its flow and heat transfer. The heat transfer is more for large volume fraction of the solid. The heat transfer is less due to the decomposition of metal particle

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**Keywords:** Cu-H<sub>2</sub>O; Nano - Fluid; Viscosity; Vertical Plate; Heat Transfer; FEM.

### Nomenclature

$B_0$	Constant applied magnetic field ( $\text{Wb m}^{-2}$ )
$c_p$	Specific heat at constant pressure ( $\text{J kg}^{-1} \text{K}^{-1}$ )
$g$	Gravity acceleration ( $\text{m s}^{-2}$ )
$J$	Current density
$M$	Dimensionless magnetic field parameter
$n$	Dimensionless frequency
$Nu$	Local Nusselt number

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Pr	Prandtl number
Q	Dimensional heat source ( $\text{kJ s}^{-1}$ )
$Q_H$	Dimensionless heat source parameter ( $\text{kJ s}^{-1}$ )
S	Dimensionless suction parameter
t	Dimensionless time (s)
$T$	Local temperature of the Nano-fluid (K)
$T_w$	Wall temperature (K)
$T_\infty$	Temperature of the ambient Nano-fluid (K)
u, w	Dimensionless velocity components ( $\text{m s}^{-1}$ )
$U_0$	Characteristic velocity ( $\text{m s}^{-1}$ )
k	Thermal conductivity
$R_a$	Radiation parameter

**Greek symbols:**

$\alpha$	Thermal diffusivity ( $\text{m}^2 \text{s}^{-1}$ )
$\beta_T$	Thermal expansion coefficient ( $\text{K}^{-1}$ )
$\varepsilon$	Dimensionless small quantity ( $\ll 1$ )
$\phi$	Solid volume fraction of the Nano-particles
$\mu$	Dynamic viscosity (Pa s)
$\psi$	Kinematic viscosity ( $\text{m}^2 \text{s}^{-1}$ )
$\theta$	Dimensionless temperature
$\sigma$	Electrical conductivity ( $\text{m}^2 \text{s}^{-1}$ )
$\sigma_1$	Stefan-Boltzmann constant
$\delta$	Mean absorption Coefficient
$\gamma$	Inclination angle of the plate
$\rho$	Density

**Superscript:**

–	Dimensional quantities
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**Subscripts:**

f	Fluid
s	Solid
nf	Nano-Fluid

## 1. INTRODUCTION:

The study of convective Nano fluid heat transfer dominated the heat transfer through fluids, due to strong physical properties of the mater (solid + base fluid). Literature survey proved that the study of Nano fluid heat transfer was worth enough to implement in industries. Sarit Kumar Das and Xiang- Qi Wang stressed the lack of heat transfer mechanisms through Nano fluids and the interaction of metal particle with solvent. However, at later stages many authors attempted to study the thermal conductivity, viscosity and physical parameters. Tiwari Arun Kumar studied the thermal conductivity and viscosity of Nano fluids with various models and concluded that existing literature was still unclear to match the experimental and theoretical results.

Recently Gbadeyan, J.A reports that the increase in Brownian motion reduced the heat transfer rate. M.A.A.Hamad concluded that the Nano particle inclusion in the base fluid changes the flow pattern significantly. To match the experimental and theoretical results Hassan A.M. and GVPN Srikanth are attempted by including the chemical reaction coefficients in the diffusion equation.

In the view of above literature survey, we want to investigate the heat transfer through Nano fluid (cu-water) with a focus on decomposition of the Nano particle (cu) in water, due to the chemical reaction by considering the viscosity proposed by Graham A.L and thermal conductivity proposed by Jang and Choi.

## 2. MATHEMATICAL FORMULATION:

The unsteady three dimensional free convection flow of a Nano-fluid past a inclined permeable, semi-infinite oscillating flat plate in the presence of an applied magnetic field with constant heat source and radiation is considered. As per cartesian coordinate system  $(\bar{x}, \bar{y}, \bar{z})$ , the flow is assumed to be in the  $\bar{x}$  direction, which is taken along the plate, and  $\bar{z}$  - axis is normal to the plate. We assume that the plate has an oscillatory movement on time  $\bar{t}$  and frequency  $\bar{n}$  with the velocity  $u(0,t)$ , which is given as  $u(0,t) = U_0 [1 + \varepsilon \cos(nt)]$  where  $\varepsilon$  is a small constant parameter ( $\varepsilon \ll 1$ ) and  $U_0$  is the characteristic velocity. We consider that initially ( $t < 0$ ) the fluid and the plate are at rest. A uniform external magnetic field  $B_0$  is taken to be acting along the  $\bar{z}$ -axis. Also assume that the induced magnetic field is small compared to the external magnetic field  $B_0$ . The surface temperature is assumed to have the constant value  $T_w$  while the ambient temperature has the constant value  $T_\infty$ , where  $T_w > T_\infty$ . The conservation equation of current density  $\nabla \cdot J = 0$  gives  $J = \text{constant}$ . Since the plate is electrically non-conducting, this constant is zero. It is assumed that the plate is infinite in extent and hence all physical quantities do not depend on  $\bar{x}$  and  $\bar{y}$  but depend only on  $\bar{z}$  and  $\bar{t}$ ,

$$\text{i.e.,} \quad \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

It is further assumed that the regular fluid and the suspended Nano-particles are in thermal equilibrium and no slip occurs between them. Under Bossinesq and boundary layer approximations, the boundary layer equations governing the flow and temperature are:

$$\frac{\partial w}{\partial z} = 0 \quad (1)$$

$$\frac{\partial u}{\partial t} + w \frac{\partial u}{\partial z} = \frac{1}{\rho_{nf}} \left[ \mu_{nf} \frac{\partial^2 u}{\partial z^2} + (\rho \beta_T)_{nf} g (T - T_\infty) \cos \gamma - \sigma B_0^2 u \right] \quad (2)$$

$$\frac{\partial T}{\partial t} + w \frac{\partial T}{\partial z} = \alpha_{nf} \frac{\partial^2 T}{\partial z^2} - \frac{Q}{(\rho c_p)_{nf}} (T - T_\infty) - \frac{1}{(\rho c_p)_{nf}} \frac{\partial q_r}{\partial z} \quad (3)$$

The appropriate initial and boundary conditions for the problem are given by

$$\left. \begin{aligned} u(z, t) &= 0, T = T_{\infty} \quad \text{for } t < 0 \quad \forall z \\ u(0, t) &= U_0 \left[ 1 + \frac{\varepsilon}{2} (e^{i n t} + e^{-i n t}) \right], T(0, t) = T_w \\ u(\infty, t) &\rightarrow 0, T(\infty, t) \rightarrow T_{\infty}, \quad \varepsilon < 1 \end{aligned} \right\} t \geq 0 \quad (4)$$

Thermo-Physical properties are related as follows:

$$\rho_{nf} = (1-\phi)\rho_f + \phi\rho_s, \quad \alpha_{nf} = \frac{k_{nf}}{(\rho c_p)_{nf}}$$

$$(\rho c_p)_{nf} = (1-\phi)(\rho c_p)_f + \phi(\rho c_p)_s$$

$$(\rho\beta)_{nf} = (1-\phi)(\rho\beta)_f + \phi(\rho\beta)_s$$

$$\frac{\mu_{nf}}{\mu_f} = 1 + 2.5\phi + 4.5 \left[ \frac{1}{\frac{h}{d_p} \left( 2 + \frac{h}{d_p} \right) \left( 1 + \frac{h}{d_p} \right)^2} \right]$$

$$k_{nf} = k_f(1-\phi) + \beta_1 k_p \phi + c_1 \frac{d_f}{d_p} k_p \text{Re}^2 d_p \text{pr} \phi$$

Due to chemical Reaction by the solvent and the metal particle the rate of change of the size of the particle can be considered as:

$$\frac{dd_p}{dt} = -K d_p$$

$$\Rightarrow d_p = d_0 e^{-K t}$$

Where  $d_p$  the diameter of the metal particle,  $d_0$  is the initial size of the particle (assumed as 50 nm) and K is the reaction rate:

$$\frac{\mu_{nf}}{\mu_f} = 1 + 2.5\phi + 4.5 \left[ \frac{1}{\left( \frac{h e^{Kt}}{50} \right) \left( 2 + \frac{h e^{Kt}}{50} \right) \left( 1 + \frac{h e^{Kt}}{50} \right)^2} \right]$$

$$k_{nf} = k_f(1-\phi) + \beta_1 k_p \phi + \frac{c_1}{50} d_f e^{Kt} k_p Re^2 d_p pr \phi \quad (5)$$

where  $\beta_1 = 0.01$  is a constant for considering the kapitza resistance per unit area

$c_1 = 18 \times 10^6$  is a proportionality constant

$$Re d_p = \frac{d_p}{\nu_f} \frac{\kappa T}{3\pi \mu_f d_p l_f} = \frac{1.381 \times 10^{23} T}{\nu_f 3\pi \mu_f (0.738)}$$

$d_f = 0.384$  nm for water

$$Pr = \text{Prandtl number} = \frac{\nu_f}{\alpha_f}$$

$l_f = \text{Mean free path} = 0.738$

$\kappa = \text{Boltzman n constant}, T = 300 K$

The thermo-physical properties (values) of the materials used are as follows:

Table		
Physical Properties	Water	Copper(Cu)
$C_p (J/kg K)$	4,179	385
$\rho (kg/m^3)$	997.1	8,933
$\kappa (W/m K)$	0.613	400
$\beta_T \times 10^{-5} (1/K)$	21	1.67
$\mu$	$8.94 \times 10^{-4}$	-----

We consider the solution of Esq. (1) as  $w = -w_0$  (6)

Where the constant  $w_0$  represents the normal velocity at the plate which is positive for suction ( $w_0 > 0$ ). Thus, we introduce the following dimensionless variables:

$$z = \left( \frac{\nu_f}{U_0} \right) Z, \quad t = \left( \frac{\nu_f}{U_0^2} \right) \tau, \quad n = \left( \frac{U_0^2}{\nu_f} \right) \eta, \quad u = U U_0, \quad \theta = \frac{T - T_\infty}{T_w - T_\infty}, \quad q_r = -\frac{4\sigma_1}{3\delta} \frac{\partial T^4}{\partial y}$$

We assume that the temperature differences within the flow are sufficiently small so that the  $T^4$  can be expressed as a linear function after using Taylor series to expand  $T^4$  about the free stream temperature  $T_\infty$  and neglecting higher-order terms. This result is the following approximation:

$$T^4 \cong 4T_\infty^3 T - 3T_\infty^4$$

By using above, we obtain

$$\frac{\partial q_r}{\partial z} = -\frac{16\sigma_1}{3\delta} \frac{\partial^2 T^4 T_\infty^3}{\partial z^2} \quad (7)$$

Using equations 4,5,6,7 the Equations 2 and 3 can be written in the following dimensionless form:

$$\begin{aligned} \left[ 1 - \phi + \phi \left( \frac{\rho_s}{\rho_f} \right) \right] \left( \frac{\partial U}{\partial \tau} - S \frac{\partial U}{\partial Z} \right) &= 1 + 2.5 \phi + 4.5 \left[ \frac{1}{\left( \frac{h e K t}{50} \right) \left( 2 + \frac{h e K t}{50} \right) \left( 1 + \frac{h e K t}{50} \right)^2} \right] \frac{\partial^2 U}{\partial Z^2} \\ &+ \left[ 1 - \phi + \phi \frac{(\rho \beta_T)_s}{(\rho \beta_T)_f} \right] Gr \theta \cos \gamma - M U \\ \left[ 1 - \phi + \phi \frac{(\rho c_p)_s}{(\rho c_p)_f} \right] \left( \frac{\partial \theta}{\partial \tau} - S \frac{\partial \theta}{\partial Z} \right) &= \frac{1}{Pr} \left[ 1 - \phi + 0.01 \phi \frac{k_p}{k_f} + \frac{k_p}{k_f} \frac{\phi}{2} \frac{\rho_f^2 c_{pf} e^{Kt}}{50 \mu_f^3} 28632.9991 \times 10^{-52} \right] \frac{\partial^2 \theta}{\partial Z^2} \\ &- \frac{1}{Pr} Q_H \theta + \frac{1}{Pr} \frac{4}{3} \frac{1}{Ra} \frac{\partial^2 \theta}{\partial Z^2} \end{aligned}$$

Where the corresponding boundary conditions (4) can be written in the dimensionless form as:

$$\begin{aligned} U(z, t) = 0, \theta(z, t) = 0 \quad \text{for } t < 0 \quad \forall z \\ U(0, t) = U_0 \left[ 1 + \frac{\varepsilon}{2} (e^{int} + e^{-int}) \right], \theta(0, t) = 1 \\ U(\infty, t) \rightarrow 0, \theta(\infty, t) \rightarrow 0 \end{aligned} \quad \left. \vphantom{\begin{aligned} U(0, t) = U_0 \left[ 1 + \frac{\varepsilon}{2} (e^{int} + e^{-int}) \right], \theta(0, t) = 1 \\ U(\infty, t) \rightarrow 0, \theta(\infty, t) \rightarrow 0 \end{aligned}} \right\} \forall t \geq 0$$

Here  $Pr$  is the Prandtl number,  $S$  is the suction ( $S > 0$ ) parameter,  $M$  is the magnetic parameter,  $Ra$  is the Radiation parameter and  $Q_H$  is the heat source parameter,  $Gr$  is the Grashof number, which are defined as:

$$Pr = \frac{\nu_f}{\alpha_f}, \quad S = \frac{w_0}{U_0}, \quad M = \frac{\sigma B_0^2 \nu_f}{\rho_f U_0^2}, \quad Ra = \frac{4 \alpha \sigma_1 T_\infty^3}{\delta k_{nf}}, \quad Q_H = \frac{Q \nu_f^2}{k_f U_0^2}, \quad Gr = g \beta_{Tf} (T_w - T_\infty) \nu_f$$

Where the velocity characteristic  $U_0$  is defined as:

$$U_0 = [g \beta_{Tf} (T_w - T_\infty) \nu_f]^{1/3}$$

The local Nusselt number (Nu) in dimensionless form:

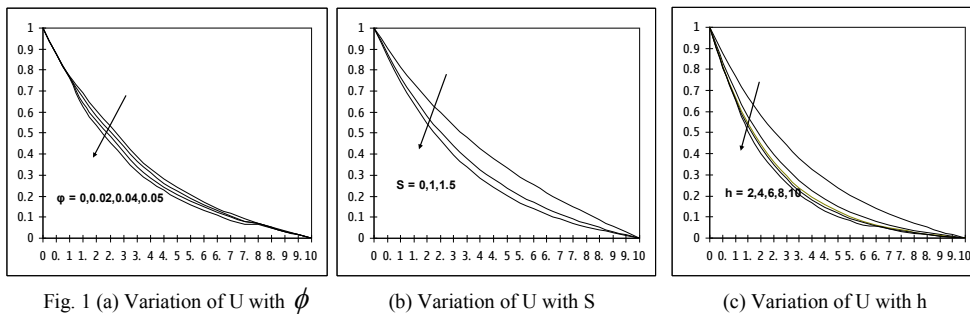
$$\text{Nu} = - \frac{k_{nf}}{k_f} \theta'(0)$$

### 3. SOLUTION OF THE PROBLEM:

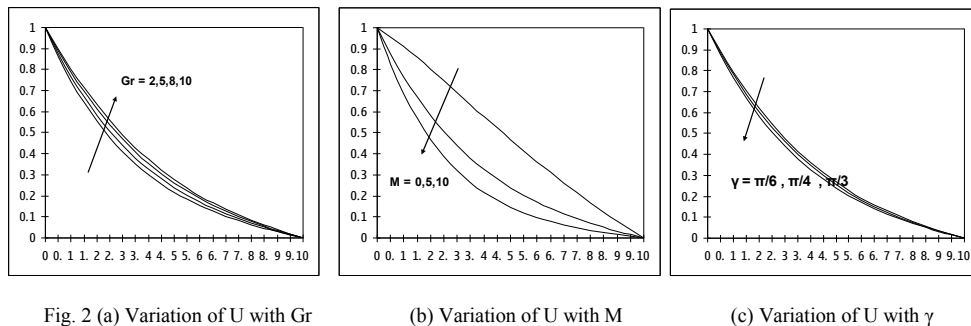
The governing equations are solved by using Galerkin Finite Element Method (FEM) as given in J. N. Reddy. The one dimensional study along the plate with three noded elements are considered and the entire domain is divided into 100 elements, for computational purpose the infinite plate is limited to 6 units. We have chosen  $\text{Pr} = 6.2$ ,  $\text{nt} = \pi/2$ ,  $n = 10$ ,  $\varepsilon = 0.02$ .

### 4. RESULTS:

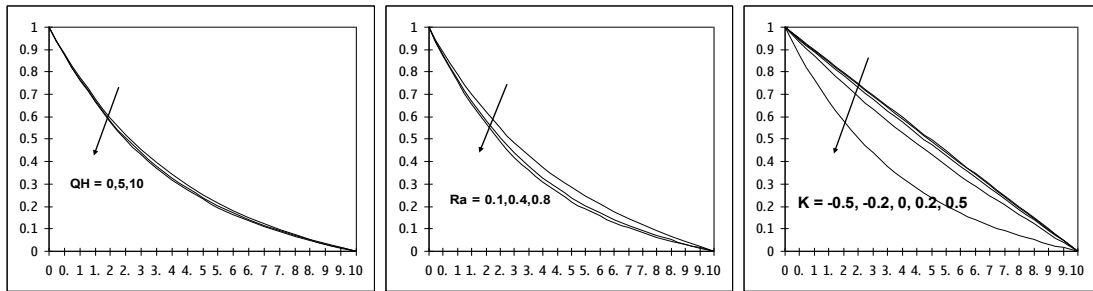
The momentum decreases with increase in volume fraction ( $\phi$ ) from Fig.1 (a). The 5% of solid volume fraction significantly decreases the momentum when compared with usual fluid. From Fig. 1(b) the suction (S) affects the momentum.



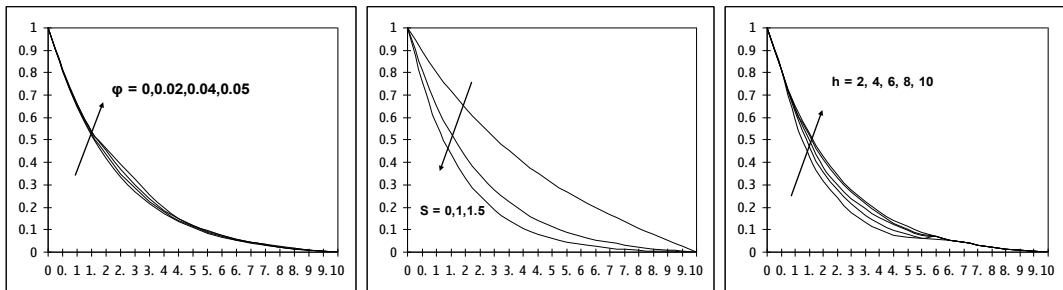
As the suction increases, the momentum decreases more rapidly compare with no suction. The increase in  $h$  drastically drops the momentum from Fig. 1(c). This occurs due to the adhesion of the layer around the Nano-particle or more packing fraction of the particles. From Fig. 2(a) the increase in thermal buoyancy the momentum increases unlike the other profiles of the momentum.



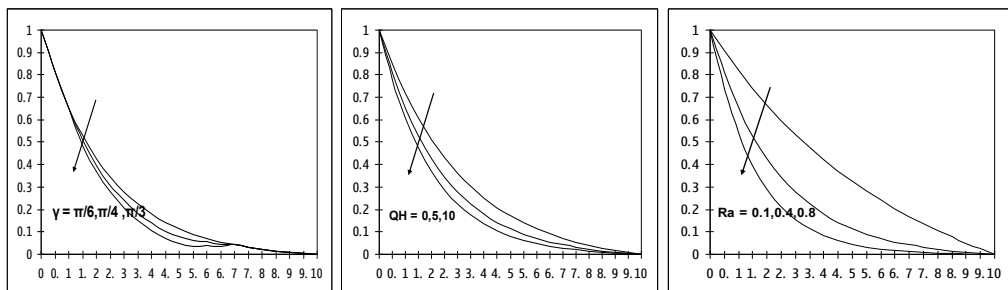
From Fig. 2(b) the magnetic field ( $M$ ) affects the momentum very much. The momentum is almost linear in the absence of the magnetic field and decreases with increase in  $M$ . From Fig.2(c) the momentum decreases slightly with increase in inclination angle ( $\gamma$ ).

Fig. 3 (a) Variation of  $U$  with  $Q_H$ (b) Variation of  $U$  with  $Ra$ (c) Variation of  $U$  with  $K$ 

From Fig. 3(a) the heat source ( $Q_H$ ) has not having much impact on momentum. But, the momentum decreases with increase in heat source. From Fig. 3(b) the radiation decreases the momentum with increase in Radiation ( $Ra$ ). It is also found that there is no impact on momentum for  $Ra > 0.4$ . From Fig. 3(c) the momentum variation is observed for generative ( $K < 0$ ) and destructive ( $K > 0$ ) chemical reactions. The generative chemical reaction produces almost the uniform momentum whereas; the destructive chemical reaction tends to exponential decay in the momentum along the plate.

Fig. 4(a) Variation of  $\theta$  with  $\phi$ (b) Variation of  $\theta$  with  $S$ (c) Variation of  $\theta$  with  $h$ 

The temperature rises with increase in the solid volume fraction ( $\phi$ ) from Fig. 4(a). From Fig. 4(b) interestingly the suction ( $S$ ) reduces the temperature along the plate almost exponentially. From Fig. 4(c) the layer ( $h$ ) acts as a semimetal to enhance the temperature along the plate.

Fig. 5(a) Variation of  $\theta$  with  $\gamma$ (b) Variation of  $\theta$  with  $Q_H$ (c) Variation of  $\theta$  with  $Ra$



From Fig. 5(a) the inclination angle ( $\gamma$ ) of the plate reduces the temperature. The thermal boundary layer has thinned by layer around the nano- particle and the inclination of the plate. From Fig. 5(b) the heat source ( $Q_H$ ) decreases the temperature. From Fig. 5(c) the radiation drastically drops the temperature even for slightly higher values ( $Ra > 0.1$ ). From Fig. 6(a) the generative chemical reaction ( $K$ ) has no impact on temperature whereas, the destructive chemical reaction reduces the temperature slightly.

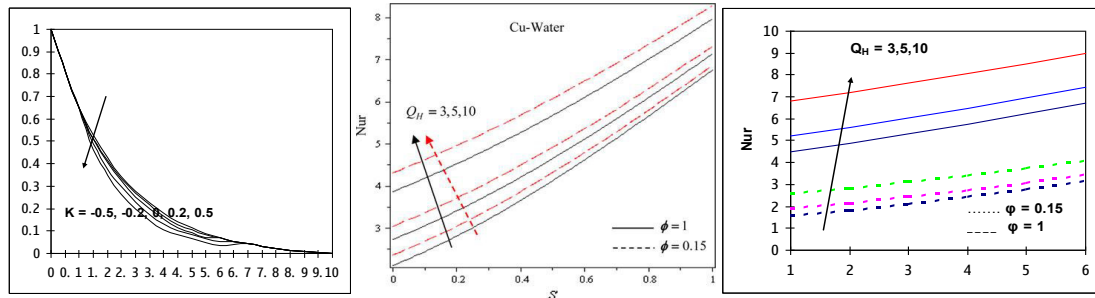


Fig. 6 (a) Variation of  $\theta$  with  $K$

(b) Variation of Nusselt Number (M.A.A.Hamad)

(c) Variation of Nusselt Number

Fig. 6(b) depicts the variation of heat transfer by M. A. A. Hamad and I.Pop. Fig. 6(c) depicts the variation of heat transfer in the present study for  $K = 0.5$  at the base of the plate. As the suction and heat source increases the heat transfer rate increases for  $\phi = 0.15, 1$ . However, the heat transfer rate has no much variation for different  $S$ . The heat transfer is more for large volume fraction of the solid. The heat transfer is less when compared with above authors due to the decomposition of metal particle.

## 5. Conclusion:

- The destructive chemical reaction exponentially decreases the velocity due to adhesion of the liquid layer around the metal particle.
- Both types of chemical reactions (destructive and generative) reduce the temperature along the plate and the variation is very narrow.

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